

# Minimal Equational Theories for Quantum Circuits

16th July 2024 - QPL'24

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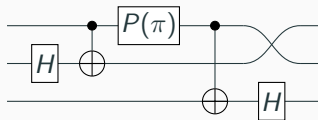
Alexandre Clément\*, Noé Delorme<sup>†</sup> and Simon Perdrix<sup>†</sup>

\* Université Paris-Saclay, ENS Paris-Saclay, CNRS, Inria, LMF, 91190, Gif-sur-Yvette, France

<sup>†</sup> Université de Lorraine, CNRS, Inria, LORIA, F-54000 Nancy, France

# What is it all about?

Quantum circuits are a rigorous graphical language used to represent quantum algorithms.

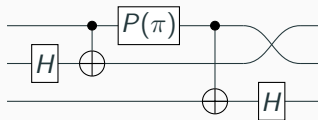


Just like boolean circuits are a rigorous graphical language used to represent classical algorithms.

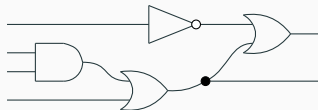


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# Quantum circuits as a graphical language

Quantum circuits are generated by the universal gateset



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to form new circuits.

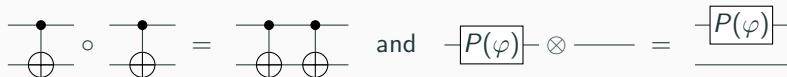


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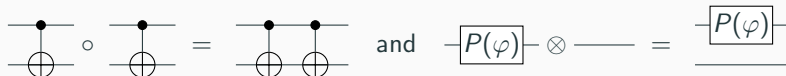


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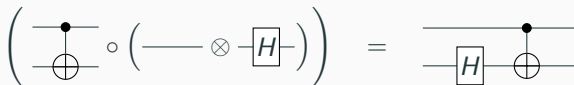
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# Standard interpretation of quantum circuits

## Interpretation

$$[[C_2 \circ C_1]] = [[C_2]] \circ [[C_1]]$$

$$[[C_1 \otimes C_2]] = [[C_1]] \otimes [[C_2]]$$

$$[[\text{I}]] = (1)$$

$$[[\text{R}(\varphi)]] = (e^{i\varphi})$$

$$[[\text{CNOT}]] = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

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## Definition “up to deformation”

Formally, quantum circuits are defined as a [symmetric monoidal category](#), which ensure some deformation equations such that

$$\boxed{P(\varphi)} \circ \text{---} = \boxed{P(\varphi)} \quad \text{or} \quad \text{---} \times \text{---} = \text{---}$$

This framework captures the intuitive behaviour of wires by ensuring that circuits are defined “up to deformation”.

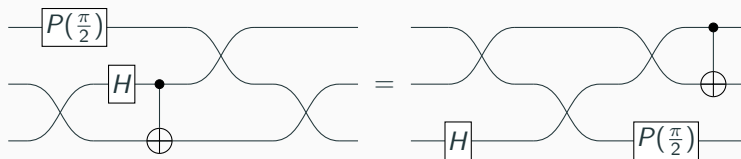


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## Other gates as shortcut notations

Other usual gates can be defined as [shortcut notation](#) by composition of the generators.

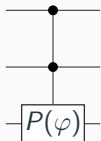
$$\boxed{Z} := \boxed{P(\pi)}$$

$$\boxed{X} := \boxed{H} \boxed{Z} \boxed{H}$$

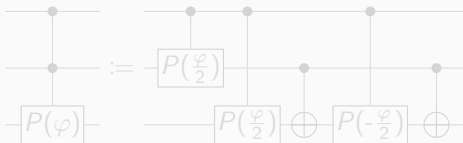
$$\boxed{R_X(\theta)} := \textcircled{-\theta/2} \boxed{H} \boxed{P(\theta)} \boxed{H}$$

## Controlled gates as shortcut notations

We use the standard [bullet notation for controlled gates](#).



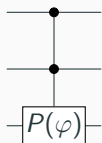
Controlled gates can be constructed [inductively](#). The  $(n + 1)$ -controlled gate is a shortcut containing several instances of  $n$ -controlled gates.



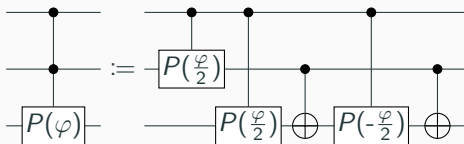
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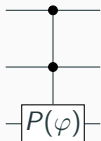
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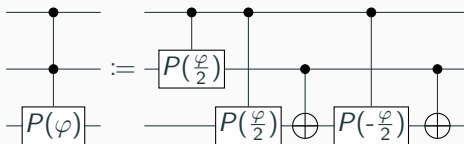
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# Motivations

Distinct circuits can have the same interpretation.

$$\left[ \begin{array}{c} \boxed{P(\frac{\pi}{2})} \\ \boxed{P(\frac{\pi}{2})} \end{array} \right] = \left[ \begin{array}{c} \boxed{H} \\ \boxed{H} \end{array} \right] = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

Given a quantum algorithm, which circuit is the best?

## Motivations:

- Resource optimisation (for instance the number of gates).
- Hardware-constraint satisfaction (for instance topological constraints).
- Verification, circuit equivalence testing.

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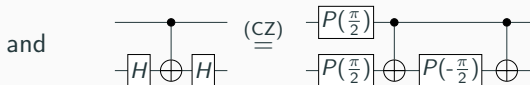
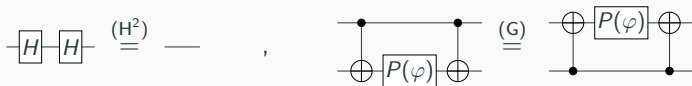
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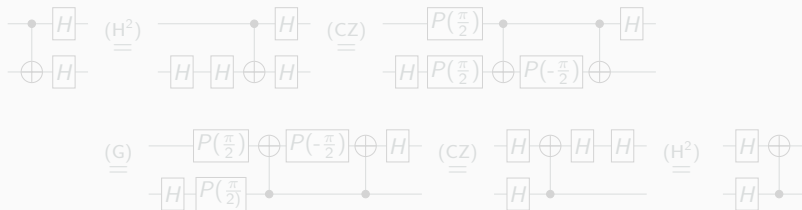


# Using equations to transform circuits

We can use simple equations such that,

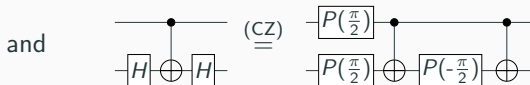
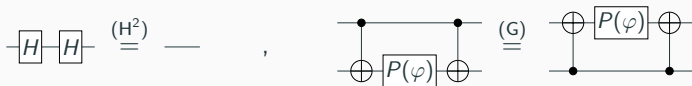


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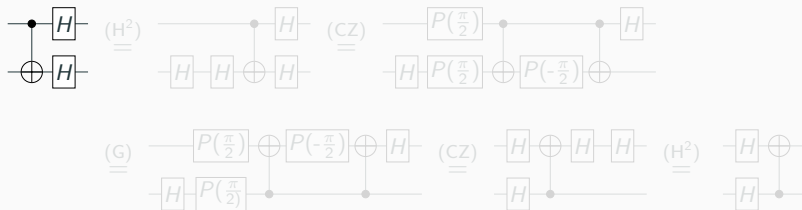


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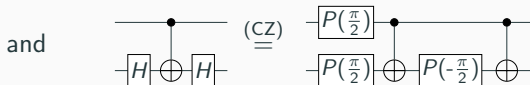
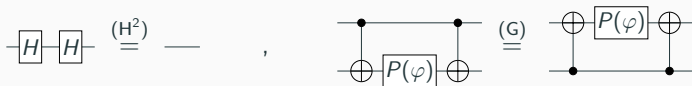


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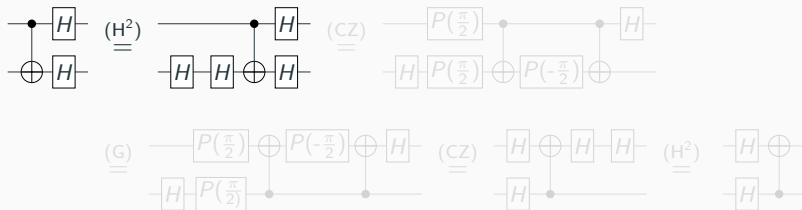


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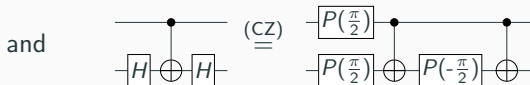
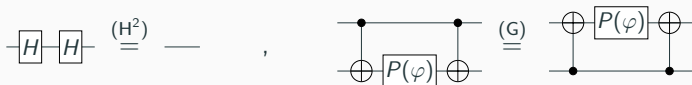


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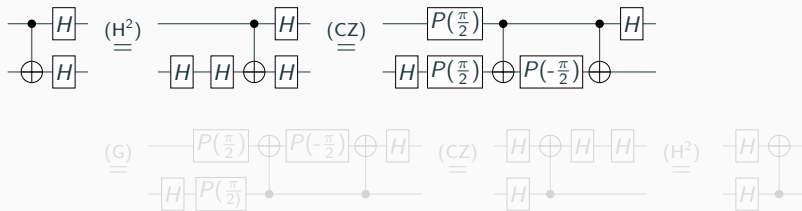


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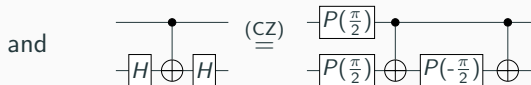
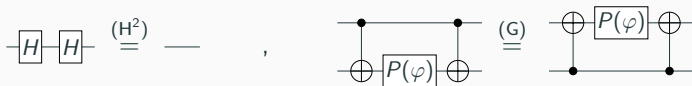


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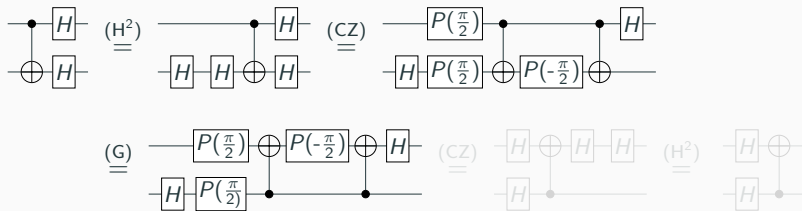


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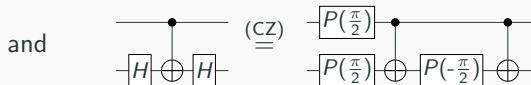
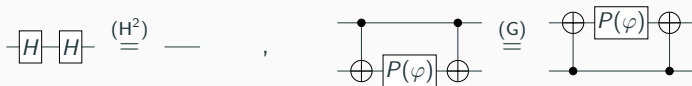


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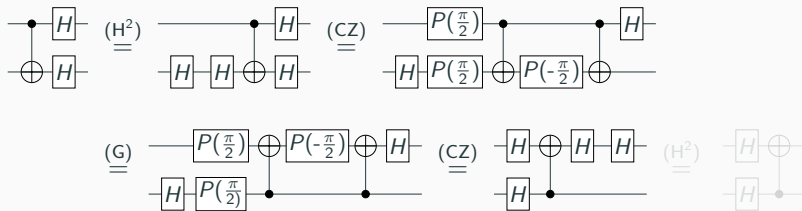


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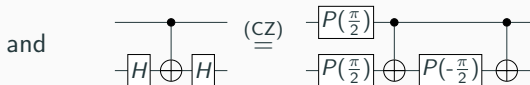
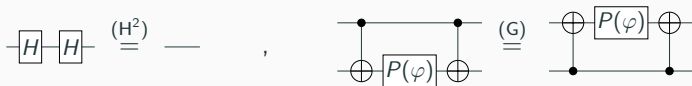


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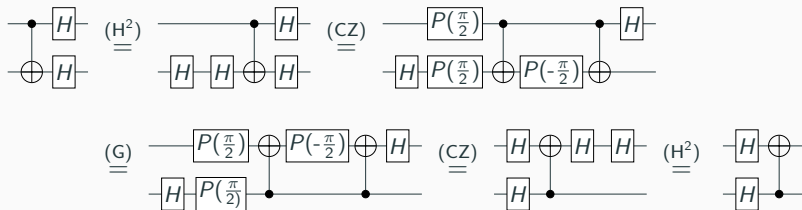


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# Soundness and completeness

Is there an **equational theory** (i.e. a set of axioms)  $\Gamma$  from which we can derive any true equation and only true equations?

## Soundness

Any derivable equation is true.

$$\forall C_1, C_2 \quad : \quad \Gamma \vdash C_1 = C_2 \implies \llbracket C_1 \rrbracket = \llbracket C_2 \rrbracket$$

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[Clément, Heurtel, Mansfield, Perdrix, Valiron'2023]

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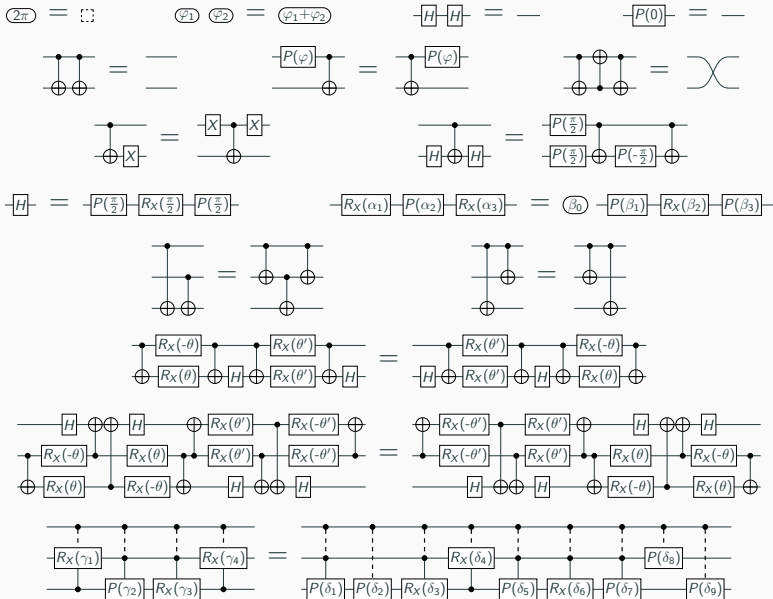
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# Complete and sound equational theory



# Euler decomposition equation

This equation follows from the well-known [Euler-decomposition](#) which states that any unitary can be decomposed, up to a global phase, into basic X- and Z-rotations.

$$\boxed{R_X(\alpha_1)} \boxed{P(\alpha_2)} \boxed{R_X(\alpha_3)} = \textcircled{\beta_0} \boxed{P(\beta_1)} \boxed{R_X(\beta_2)} \boxed{P(\beta_3)}$$

It represents a [family of equations](#): there are explicit trigonometric relations to compute  $\beta_0, \beta_1, \beta_2, \beta_3$  as functions of  $\alpha_1, \alpha_2, \alpha_3$ .

By choosing specific parameters, we can retrieve simple equations, such that

$$\boxed{P(\varphi_1)} \boxed{P(\varphi_2)} = \boxed{P(\varphi_1 + \varphi_2)} \quad \boxed{X} \boxed{P(\varphi)} \boxed{X} = \textcircled{\varphi} \boxed{P(-\varphi)}$$

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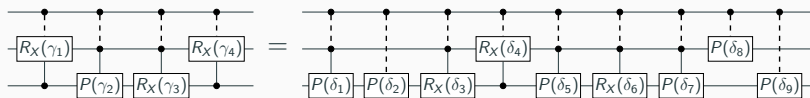
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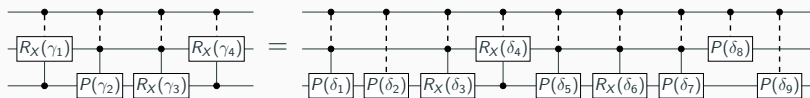
## The weird equation



Similarly to the Euler decomposition equation, it represents a **family of equations**: there is an instance of this equation in the equational theory for any number of wires  $n \geq 2$  and for any parameters  $\gamma_1, \gamma_2, \gamma_3, \gamma_4 \in \mathbb{R}$ .

The presence of such weird equation is the consequence of the technique used to prove completeness: the proof is based on **back and forth translations between quantum circuits and optical circuits**.

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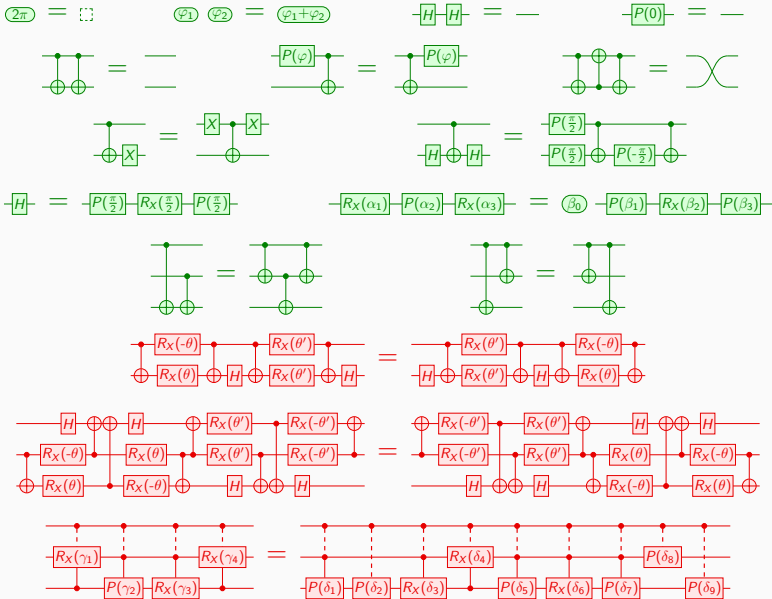


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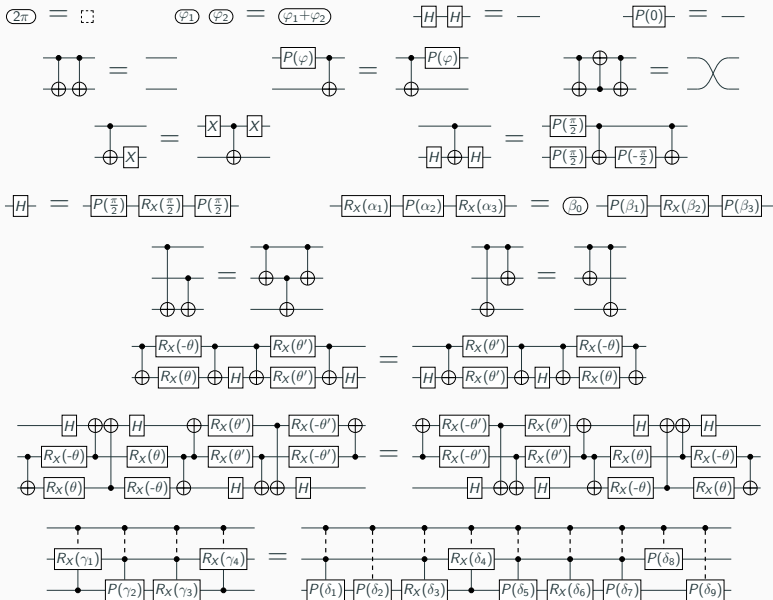
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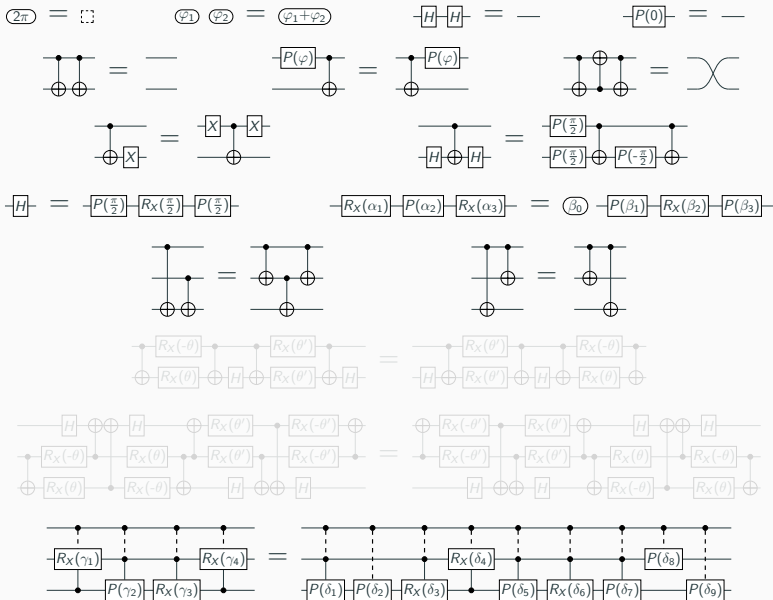
# Some **easy** and some **intricate** equations



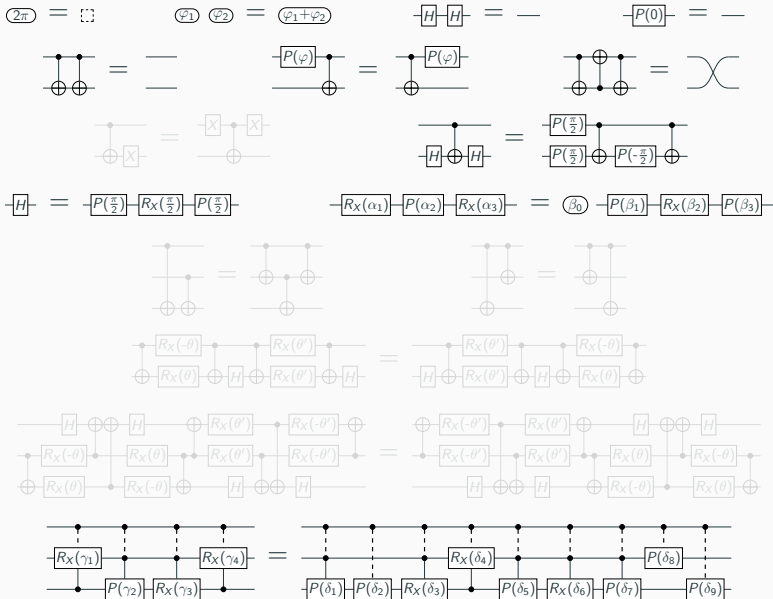
# Simplifications



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$$\boxed{2\pi} = \boxed{\square} \quad \boxed{\varphi_1} \boxed{\varphi_2} = \boxed{\varphi_1 + \varphi_2} \quad \boxed{H} \boxed{H} = \text{---} \quad \boxed{P(0)} = \text{---}$$

$$\begin{array}{c} \bullet \\ \oplus \end{array} \begin{array}{c} \boxed{P(\varphi)} \\ \oplus \end{array} = \boxed{P(\varphi)} \text{---}$$

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$$\boxed{H} = \boxed{P(\frac{\pi}{2})} \boxed{R_X(\frac{\pi}{2})} \boxed{P(\frac{\pi}{2})}$$

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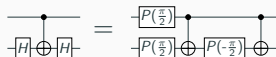
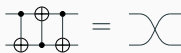
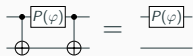
$$\begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(-\theta)} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(\theta')} \begin{array}{c} \oplus \\ \oplus \end{array} = \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(\theta')} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(-\theta)} \begin{array}{c} \oplus \\ \oplus \end{array}$$

$$\begin{array}{c} \oplus \\ \oplus \end{array} \boxed{H} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{H} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(\theta')} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(-\theta')} \begin{array}{c} \oplus \\ \oplus \end{array} = \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(-\theta')} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{R_X(\theta')} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{H} \begin{array}{c} \oplus \\ \oplus \end{array} \boxed{H} \begin{array}{c} \oplus \\ \oplus \end{array}$$

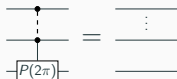
$$\begin{array}{c} \bullet \\ \oplus \end{array} \boxed{R_X(\gamma_1)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\gamma_2)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{R_X(\gamma_3)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{R_X(\gamma_4)} = \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\delta_1)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\delta_2)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{R_X(\delta_3)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\delta_5)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{R_X(\delta_6)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\delta_7)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\delta_8)} \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\delta_9)}$$

# Simplifications

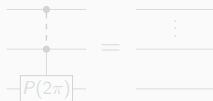
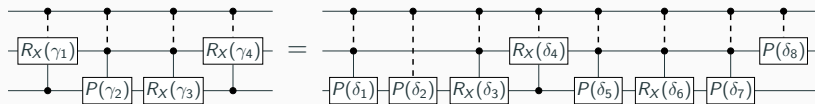
$$\boxed{2\pi} = \boxed{\quad} \quad \boxed{\varphi_1} \boxed{\varphi_2} = \boxed{\varphi_1 + \varphi_2} \quad \boxed{H} \boxed{H} = \text{---} \quad \boxed{P(0)} = \text{---}$$



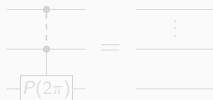
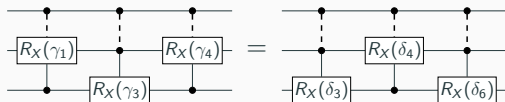
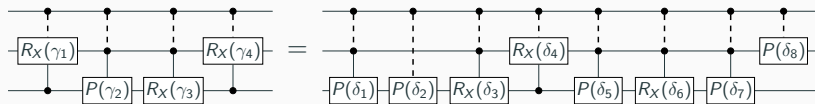
$$\boxed{H} = \boxed{P(\frac{\pi}{2})} \boxed{R_X(\frac{\pi}{2})} \boxed{P(\frac{\pi}{2})} \quad \boxed{R_X(\alpha_1)} \boxed{P(\alpha_2)} \boxed{R_X(\alpha_3)} = \boxed{\beta_0} \boxed{P(\beta_1)} \boxed{R_X(\beta_2)} \boxed{P(\beta_3)}$$



# Killing the remaining weird rule

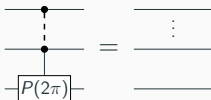
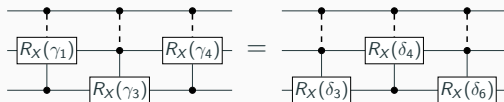
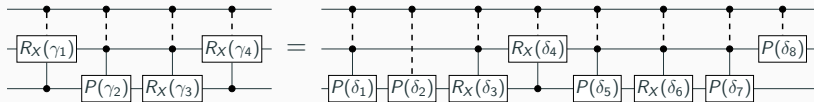


# Killing the remaining weird rule





# Killing the remaining weird rule



# Towards the limit of simplification

$$\begin{array}{l}
 \boxed{2\pi} = \boxed{\quad} \quad \boxed{\varphi_1} \boxed{\varphi_2} = \boxed{\varphi_1 + \varphi_2} \quad \boxed{H} \boxed{H} = \text{---} \quad \boxed{P(0)} = \text{---} \\
 \begin{array}{c} \bullet \\ \oplus \end{array} \boxed{P(\varphi)} \begin{array}{c} \bullet \\ \oplus \end{array} = \boxed{P(\varphi)} \text{---} \quad \begin{array}{c} \bullet \\ \oplus \end{array} \begin{array}{c} \oplus \\ \bullet \end{array} = \begin{array}{c} \text{---} \\ \oplus \end{array} \quad \begin{array}{c} \text{---} \\ \oplus \end{array} = \boxed{P(\frac{\pi}{2})} \begin{array}{c} \bullet \\ \oplus \end{array} \begin{array}{c} \text{---} \\ \oplus \end{array} \\
 \boxed{H} = \boxed{P(\frac{\pi}{2})} \boxed{R_X(\frac{\pi}{2})} \boxed{P(\frac{\pi}{2})} \quad \boxed{R_X(\alpha_1)} \boxed{P(\alpha_2)} \boxed{R_X(\alpha_3)} = \boxed{\beta_0} \boxed{P(\beta_1)} \boxed{R_X(\beta_2)} \boxed{P(\beta_3)} \\
 \begin{array}{c} \bullet \\ \vdots \\ \oplus \end{array} \boxed{P(2\pi)} = \begin{array}{c} \text{---} \\ \vdots \\ \text{---} \end{array} \} n \geq 3
 \end{array}$$

**Question:** Can we simplify the equational theory even more?

## Theorem

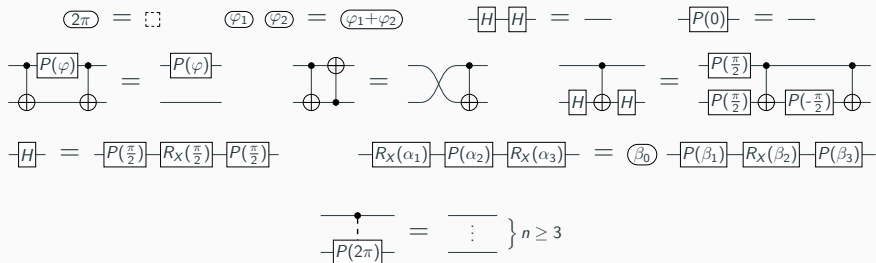
This equational theory is complete, sound and **minimal**.

## Minimality

All equations are independent.

$\forall (C_1 = C_2) \in \Gamma \quad : \quad \Gamma \setminus \{C_1 = C_2\} \not\vdash C_1 = C_2$

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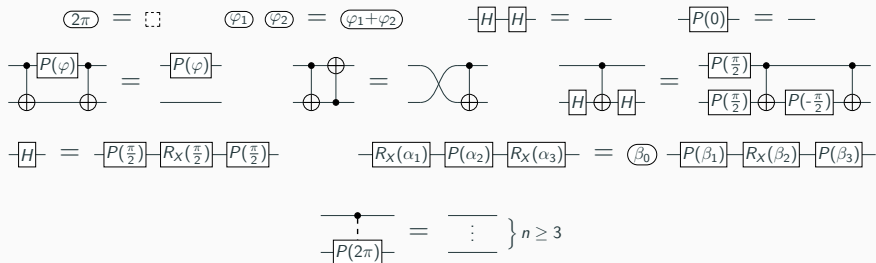
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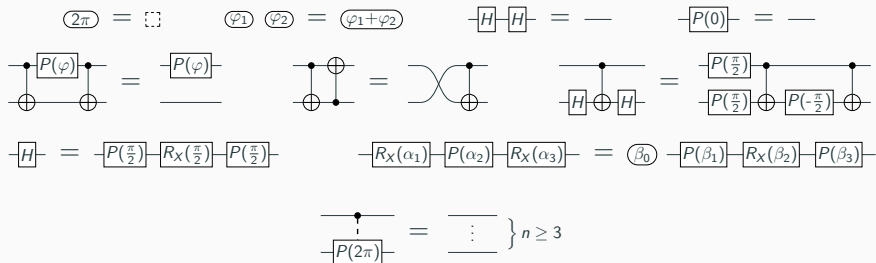
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# Necessity of the simple equations

$$\begin{array}{l}
 \textcircled{2\pi} = \boxed{\phantom{0}} \quad \textcircled{\varphi_1} \textcircled{\varphi_2} = \textcircled{\varphi_1 + \varphi_2} \quad \boxed{H} \boxed{H} = \text{---} \quad \boxed{P(0)} = \text{---} \\
 \begin{array}{c} \bullet \\ | \\ \oplus \end{array} \boxed{P(\varphi)} \begin{array}{c} \bullet \\ | \\ \oplus \end{array} = \boxed{P(\varphi)} \text{---} \quad \begin{array}{c} \bullet \\ | \\ \oplus \end{array} \begin{array}{c} \bullet \\ | \\ \oplus \end{array} = \begin{array}{c} \bullet \\ | \\ \oplus \end{array} \begin{array}{c} \bullet \\ | \\ \oplus \end{array} \\
 \boxed{H} = \boxed{P(\frac{\pi}{2})} \boxed{R_X(\frac{\pi}{2})} \boxed{P(\frac{\pi}{2})} \quad \boxed{R_X(\alpha_1)} \boxed{P(\alpha_2)} \boxed{R_X(\alpha_3)} = \textcircled{\beta_0} \boxed{P(\beta_1)} \boxed{R_X(\beta_2)} \boxed{P(\beta_3)} \\
 \begin{array}{c} \bullet \\ | \\ \vdots \\ \bullet \\ | \\ \boxed{P(2\pi)} \end{array} = \begin{array}{c} \text{---} \\ | \\ \vdots \\ \text{---} \end{array} \left. \vphantom{\begin{array}{c} \bullet \\ | \\ \vdots \\ \bullet \\ | \\ \boxed{P(2\pi)} \end{array}} \right\} n \geq 3
 \end{array}$$

For instance, the blue equation is the only one that does not preserve the parity of the number of swap gates.

# Necessity of the Euler decomposition equation

Equation (E) represent a **family of equations** and is the only equation involving **non-linear** computations.

$$\boxed{R_X(\alpha_1)} \boxed{P(\alpha_2)} \boxed{R_X(\alpha_3)} \stackrel{(E)}{=} \boxed{\beta_0} \boxed{P(\beta_1)} \boxed{R_X(\beta_2)} \boxed{P(\beta_3)}$$

Maybe (E) is in the equational theory only to retrieve simple equations such that

$$\boxed{P(\varphi_1)} \boxed{P(\varphi_2)} \stackrel{(P_+)}{=} \boxed{P(\varphi_1 + \varphi_2)} \quad \boxed{X} \boxed{P(\varphi)} \boxed{X} \stackrel{(P_-)}{=} \boxed{\varphi} \boxed{P(-\varphi)}$$

## Proposition

Let  $\Gamma$  be a set of equations containing

- all the equations of the equational theory except (E),
- any set of instance of (E) of cardinality strictly less than  $2^{\aleph_0}$ ,
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Then there exists an instance of (E) which is not a consequence of  $\Gamma$ .  
Hence, **uncountably many instances of (E) are required.**

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# Unboundedness of the equational theory

Every instances of  $\overline{\begin{array}{c} \bullet \\ | \\ \boxed{P(2\pi)} \end{array}} = \overline{\begin{array}{c} \vdots \end{array}} \}_{n \geq 3}$  are necessary (for every  $n \geq 3$ ).

## Theorem

There is **no** complete equational theory for quantum circuits made of equations acting on a **bounded number of wires**.

More precisely, any complete equational theory for quantum circuits has **at least one equation acting on  $n$  wires for any  $n \in \mathbb{N}$** .

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# Proof sketch [1/2]

## Alternative interpretation

For any  $k \in \mathbb{N}$ , for any quantum circuit  $C$ , let  $\llbracket C \rrbracket_k^\# \in [0, 2\pi)$  be inductively defined as

$$\llbracket C_2 \circ C_1 \rrbracket_k^\# = \llbracket C_1 \otimes C_2 \rrbracket_k^\# = \llbracket C_2 \rrbracket_k^\# + \llbracket C_1 \rrbracket_k^\# \bmod 2\pi$$

$$\llbracket \text{[ ]} \rrbracket_k^\# = \llbracket \text{---} \rrbracket_k^\# = 0 \quad \llbracket \text{[ } \varphi \text{ ]} \rrbracket_k^\# = 2^k \varphi \bmod 2\pi \quad \llbracket \text{[ } \text{---} \text{ ]} \rrbracket_k^\# = 2^{k-1} \pi \bmod 2\pi$$

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**Intuition:**  $\llbracket C \rrbracket_n^\# = \arg(\det(\llbracket C \rrbracket))$  for any  $n$ -qubit quantum circuit  $C$ .

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For any  $n$ -qubit quantum circuits  $C_1, C_2$  and  $k \geq n$ ,

$$\llbracket C_1 \rrbracket = \llbracket C_2 \rrbracket \implies \llbracket C_1 \rrbracket_k^\# = \llbracket C_2 \rrbracket_k^\#$$

Thus, any sound equation involving circuits acting on at most  $n - 1$  wires is also sound according to  $\llbracket \cdot \rrbracket_{n-1}^\#$ .

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Thus, any sound equation involving circuits acting on **at most  $n - 1$  wires** is also **sound** according to  $\llbracket \cdot \rrbracket_{n-1}^\#$ .

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$$\left[ \left[ \begin{array}{c} \text{---} \bullet \text{---} \\ | \\ \boxed{P(2\pi)} \\ \text{---} \end{array} \right] \right\}_n \right]_{n-1}^\# = \pi \neq 0 = \left[ \left[ \begin{array}{c} \text{---} \\ \vdots \\ \text{---} \end{array} \right] \right\}_n \right]_{n-1}^\#$$

Hence  $\left[ \begin{array}{c} \text{---} \bullet \text{---} \\ | \\ \boxed{P(2\pi)} \\ \text{---} \end{array} \right] = \left[ \begin{array}{c} \text{---} \\ \vdots \\ \text{---} \end{array} \right] \}_n$  **cannot be derived** from an equational theory containing only **equations acting on strictly less than  $n$  wires**.

## Discussion of the theorem

**Possible weakness:**  $\|C\|_k^\#$  is closely related to the determinant of  $\|C\|$ .  
What if we consider quantum circuits **up to global phases**?

→ The theorem still holds!

**Possible weakness:** The choice of the generators  $\boxed{H}$ ,  $\boxed{P(\varphi)}$ ,  $\text{CNOT}$ ,  $\text{Phase}$  may seem arbitrary. What if we take **another universal gate set**?

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# Interesting corollaries

## Corollary

Any complete equational theory for the **fragment where parameters are multiple of  $\frac{\pi}{2^n}$**  must contain at least one equation acting on  **$n + 2$  wires**.

For **Clifford** quantum circuits (case  $n = 1$ ),

→ The bound has been reached [Selinger'2015].

For **Clifford+T** quantum circuits (case  $n = 2$ ),

→ There exists equations that are **not provable** in the equational theory for 2-qubit Clifford+T of [Bian,Selinger'2022].

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# Extension to quantum circuits with ancillae

Quantum circuits **with ancillae** are generated by



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respectively denoting **qubit initialisation** and **qubit termination**.

(The generator  $\dashv$  can only be applied to qubits in the  $|0\rangle$ -state.)

## Semantics

We extend  $\llbracket \cdot \rrbracket$  with  $\llbracket \vdash \rrbracket = |0\rangle$  and  $\llbracket \dashv \rrbracket = \langle 0|$ .

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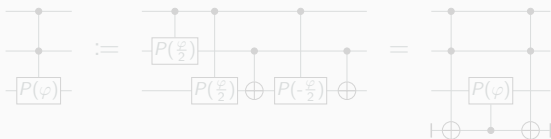
# Boundedness of the equational theory with ancillae

## Theorem [Clément, Delorme, Perdrix, Vilmart'2024]

Adding those three equations makes the equational theory **complete for quantum circuits with ancillae**.

$$\vdash \text{---} = \boxed{\text{---}} \quad , \quad \vdash \boxed{P(\varphi)} \text{---} = \text{---} \quad , \quad \begin{array}{c} \bullet \\ | \\ \text{---} \\ \oplus \\ \text{---} \end{array} = \text{---}$$

Using ancillae, we can build controlled gates **without dividing the angles**.



In these more general settings,  $\begin{array}{c} \bullet \\ | \\ \text{---} \\ \vdots \\ \text{---} \\ \oplus \\ \text{---} \\ \vdots \\ \text{---} \end{array} = \left. \begin{array}{c} \text{---} \\ \vdots \\ \text{---} \end{array} \right\} n \text{ is derivable for } n \geq 4.$

Hence, using ancillae, there is a complete equational theory **made of equations acting on at most 3 wires**.

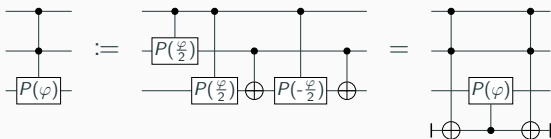
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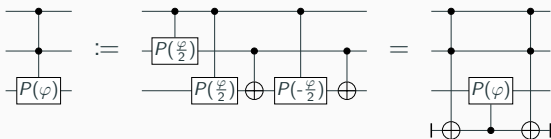
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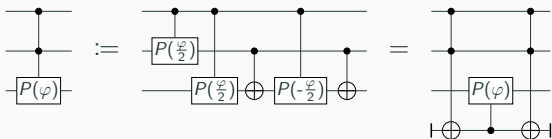
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# Thanks



arXiv:2311.07476

**Minimal Equational Theories for Quantum Circuits**

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